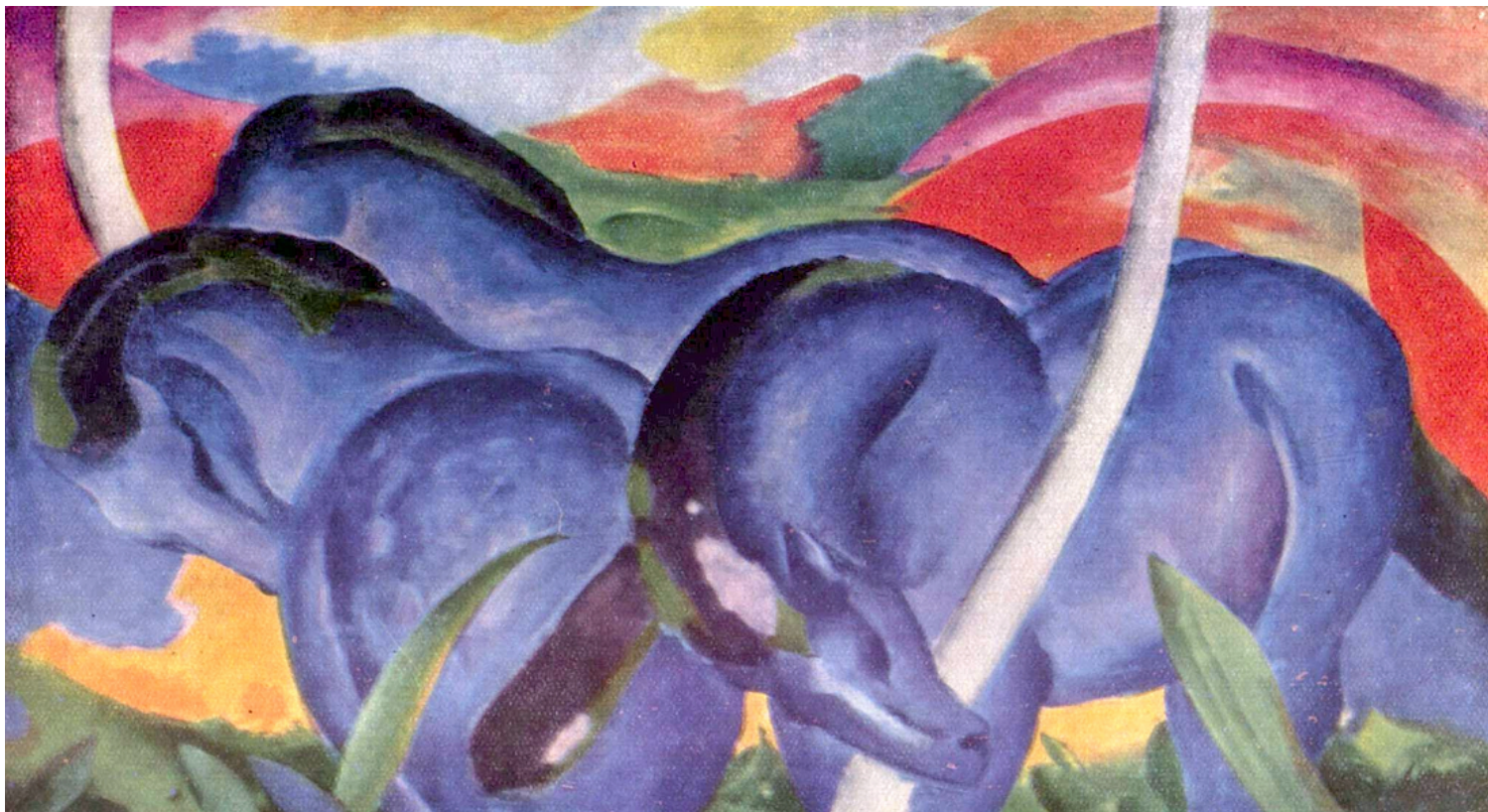


QUANTIZATION OF THREE-BODY SCATTERING AMPLITUDE

Maxim Mai

The George Washington University



Franz Marc, "Die grossen blauen Pferde" 1911

- **QCD at low energies** → *mass generation & confinement*

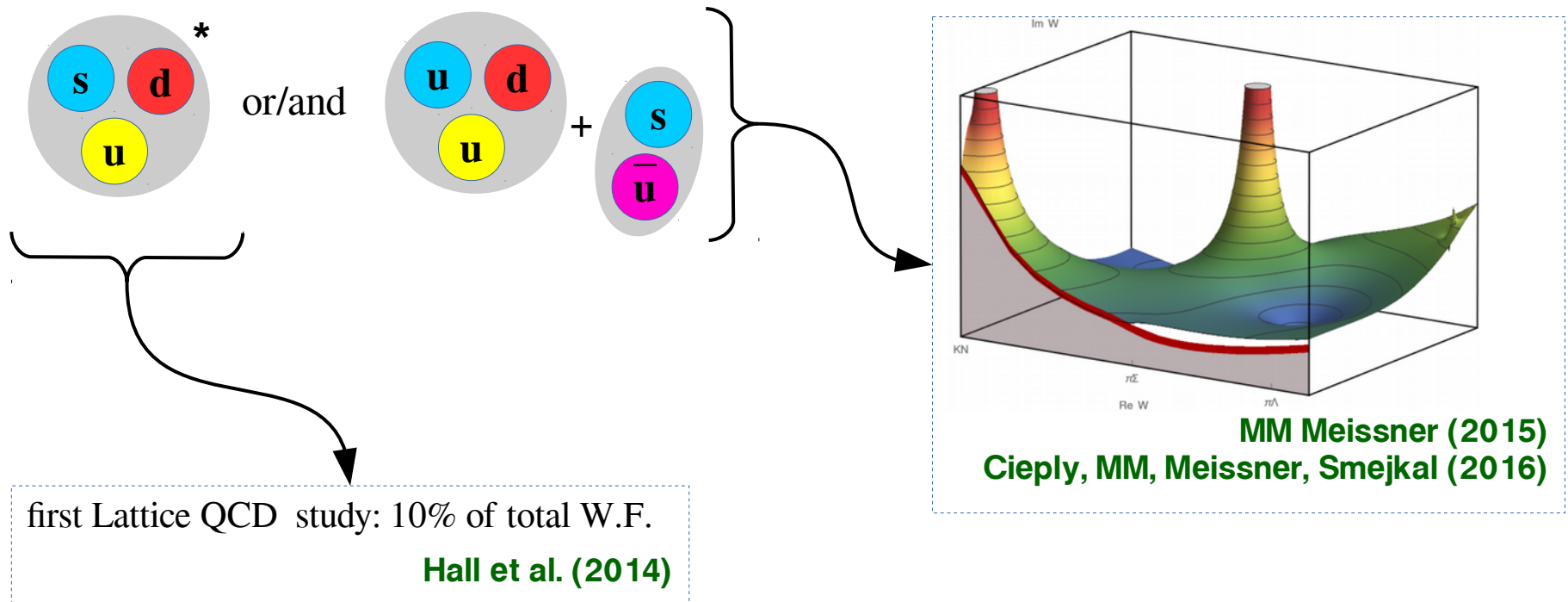
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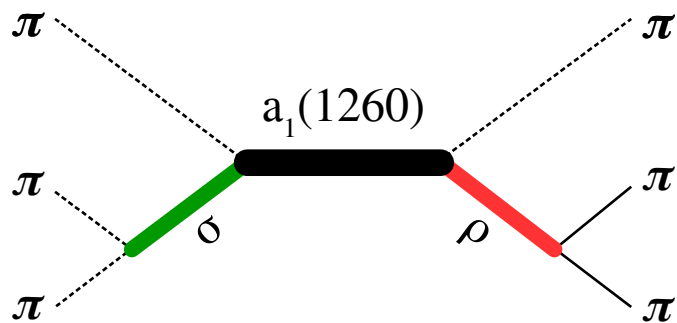
E.g.: $\Lambda(1405)$



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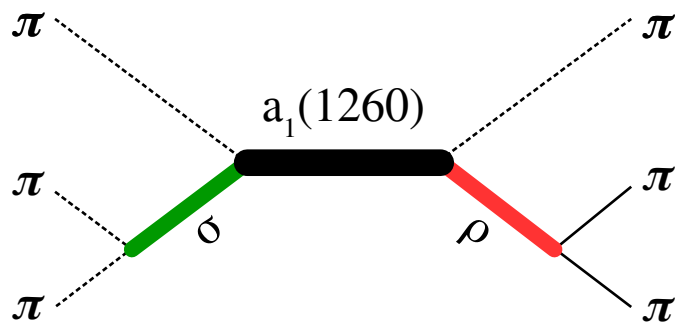


- important channel in GlueX @ JLab

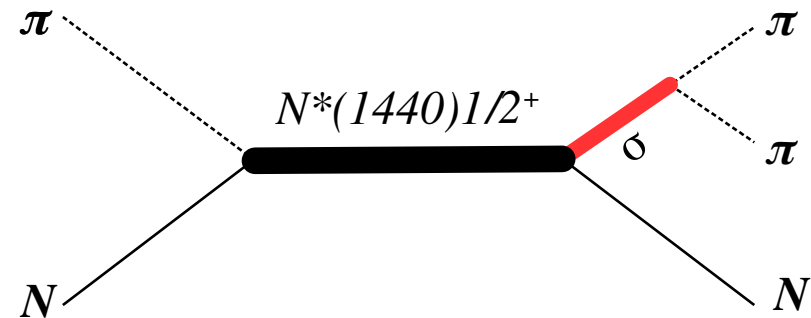


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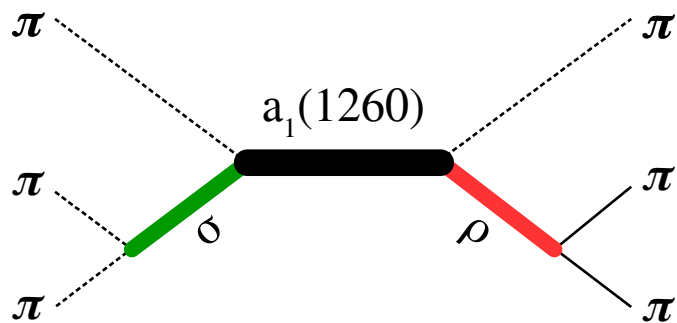
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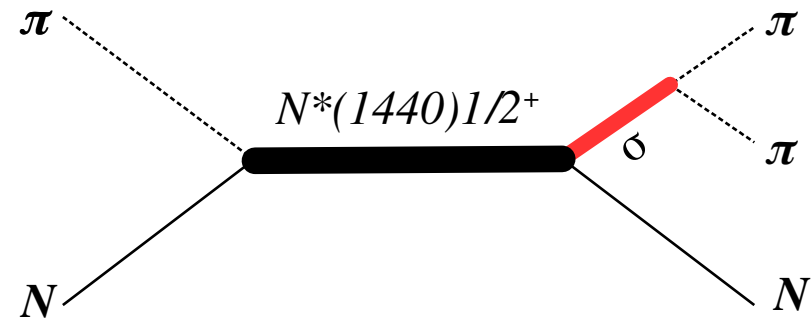
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- **first Lattice QCD results:**

w. incomplete treatment of $\pi\pi N$

→ **NO Roper-signal**

Lang et al. (2017)

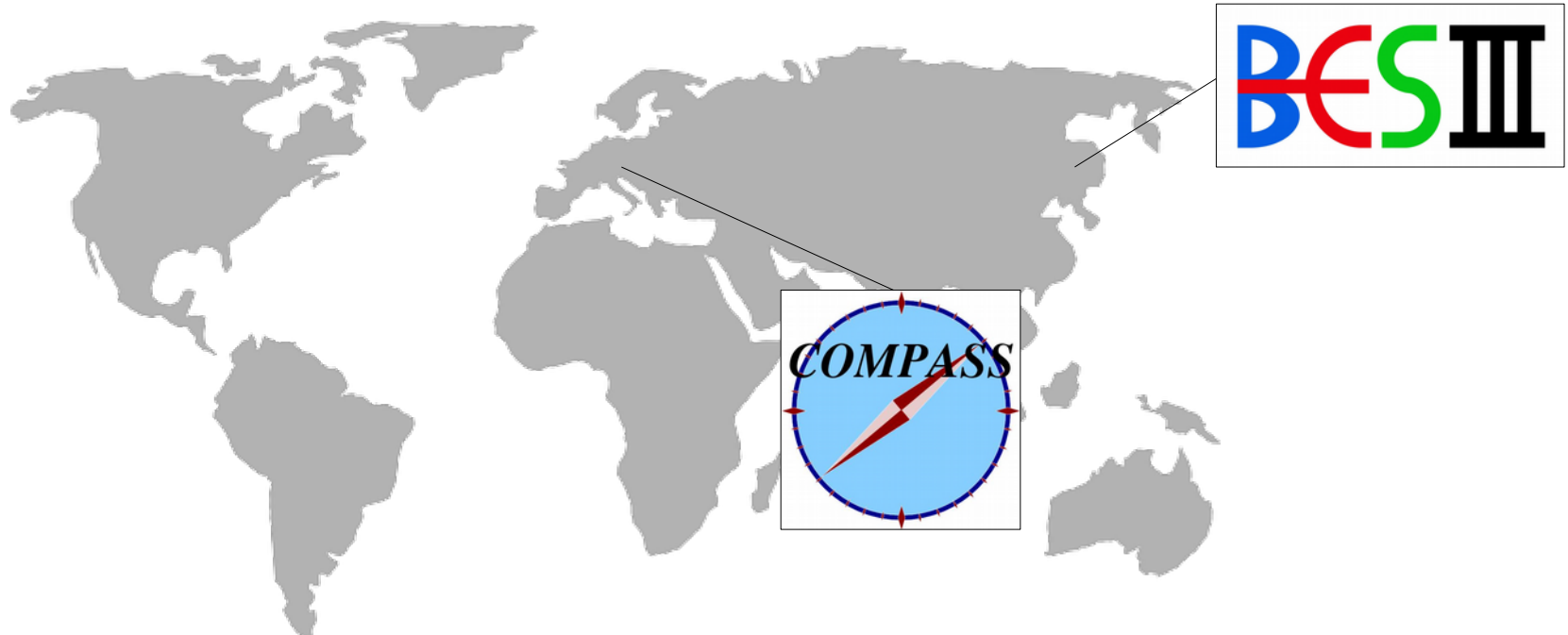
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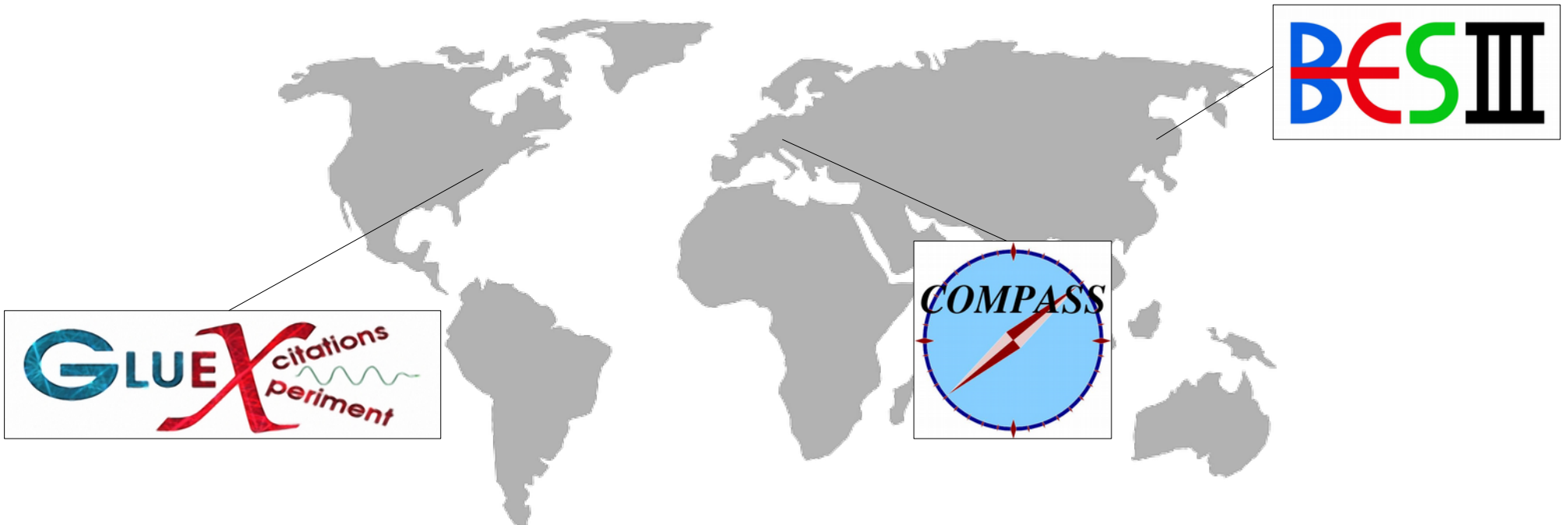
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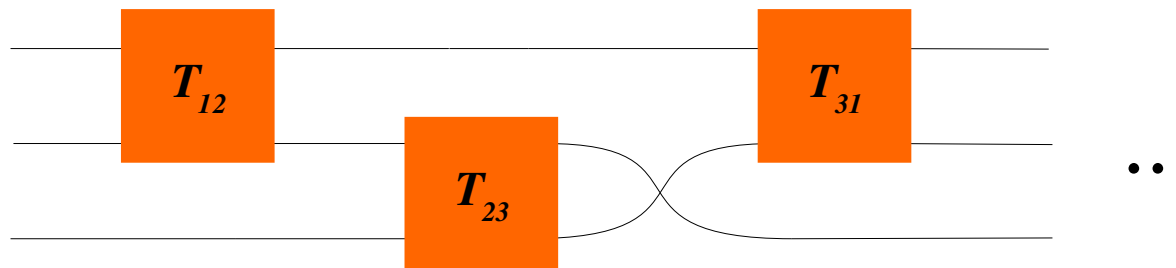
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Faddeev(1959)



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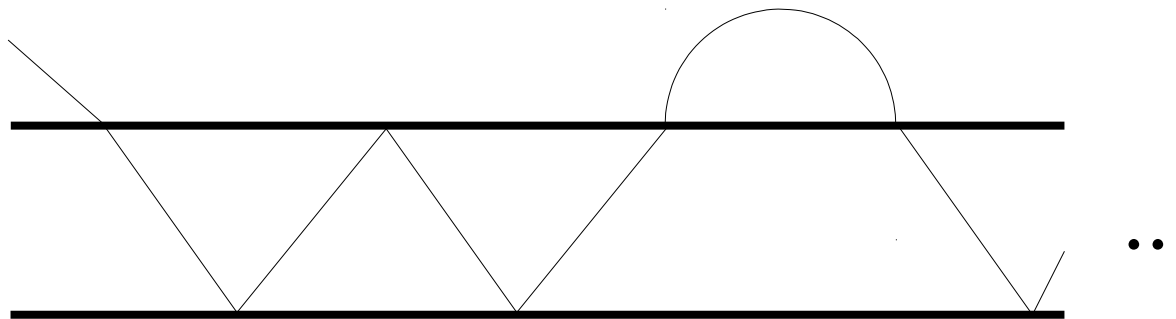
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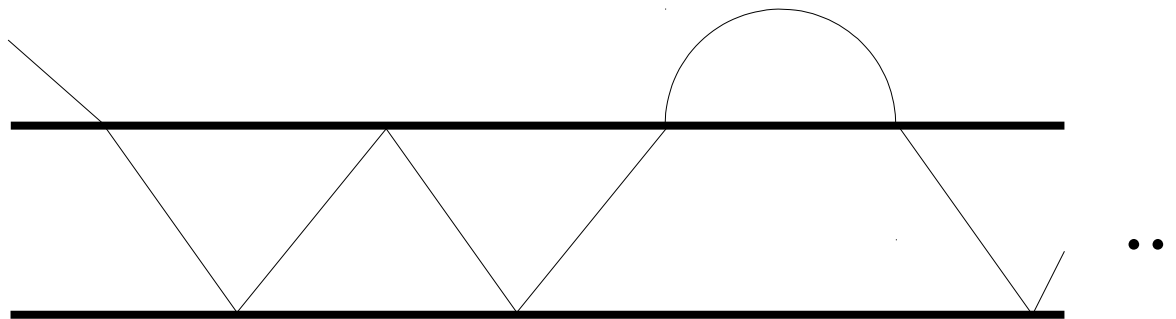
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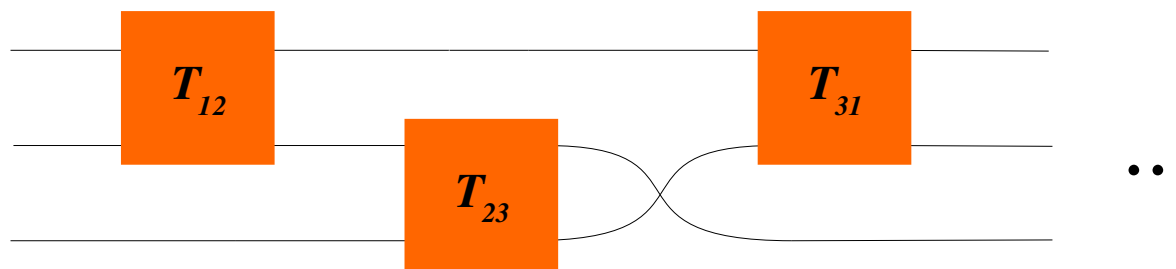
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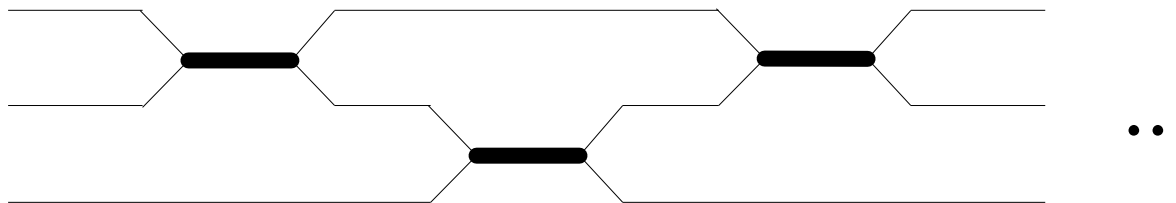


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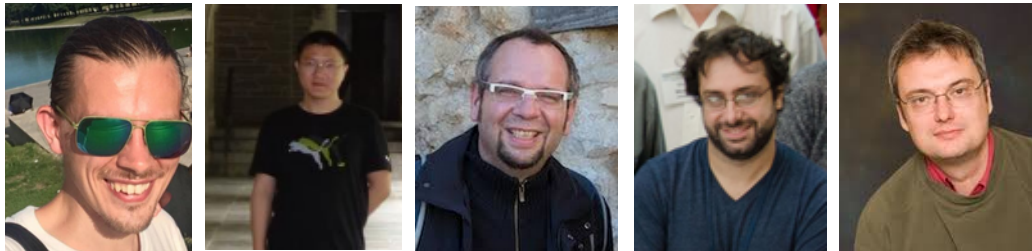
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FADDEEV EQUATIONS WITH ISOBARS

MM, Hu, Döring, Pilloni, Szczepaniak

Eur.Phys.J. A53 (2017) no.9, 177



FE in isobar parametrization

Original study – Amado Model

Amado, Aaron, Young (1968)

- 3-dimensional integral equation from unitarity constraint & BSE ansatz
- valid below break-up energies ($E < 3m$) & analyticity constraints unclear

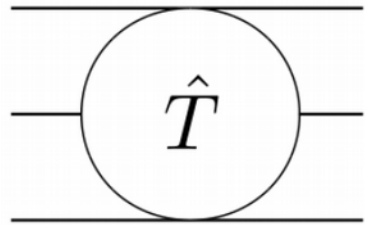
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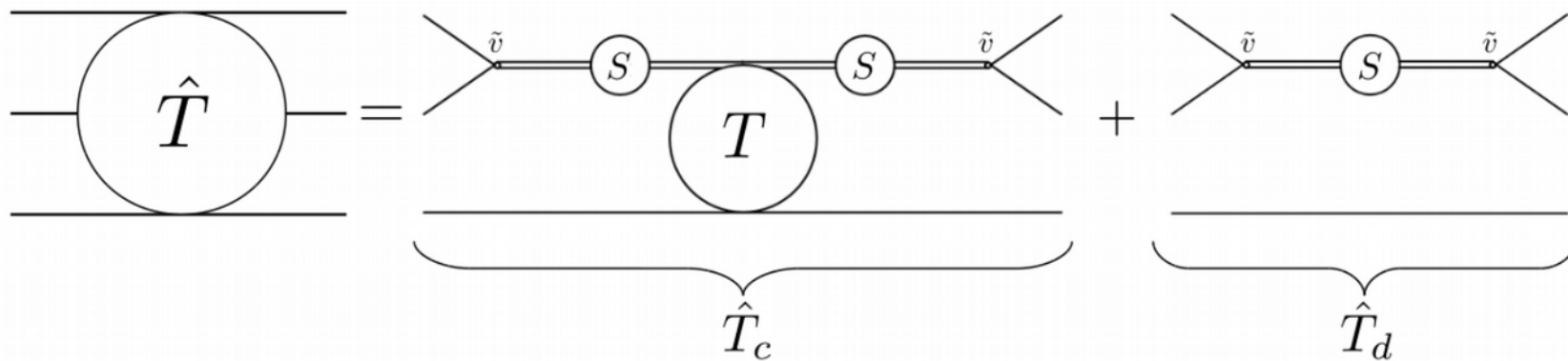
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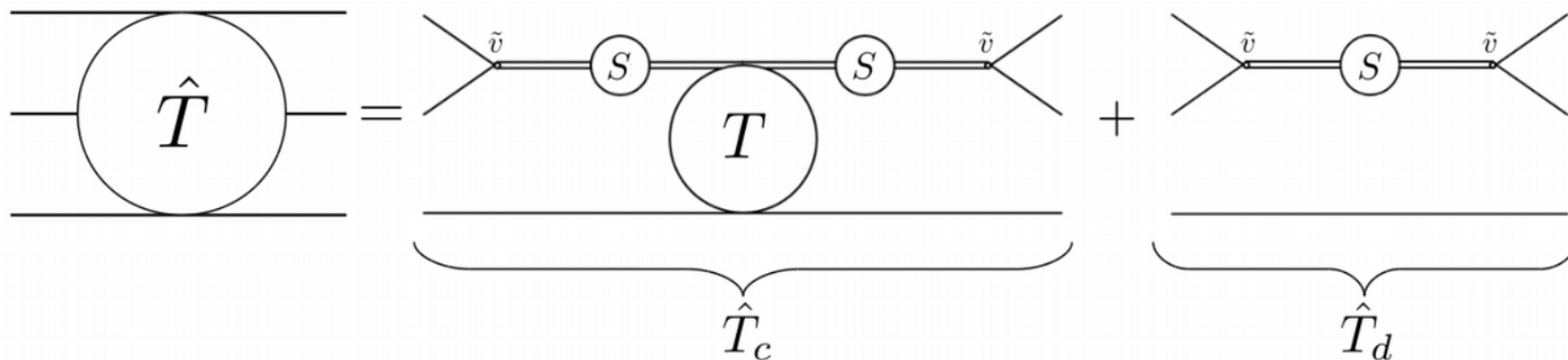
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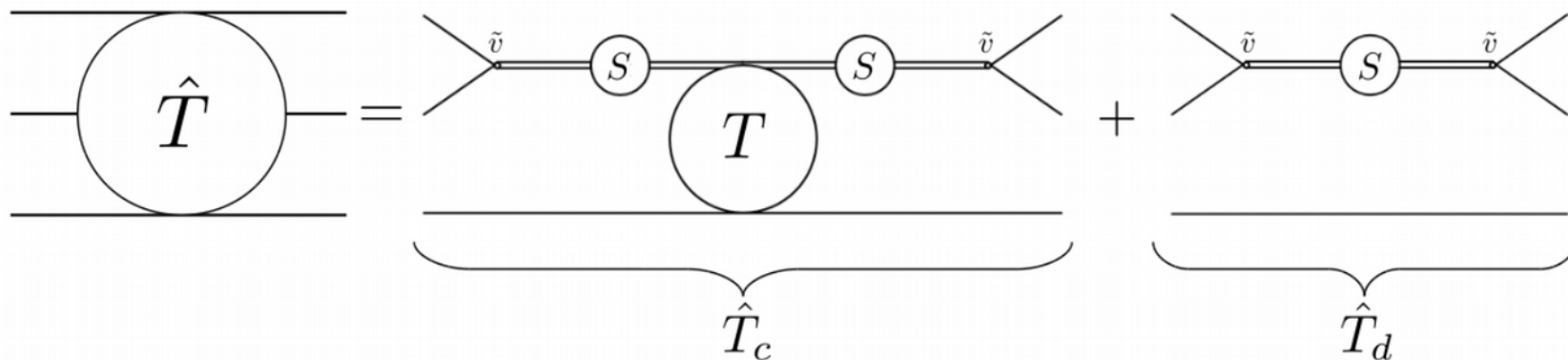
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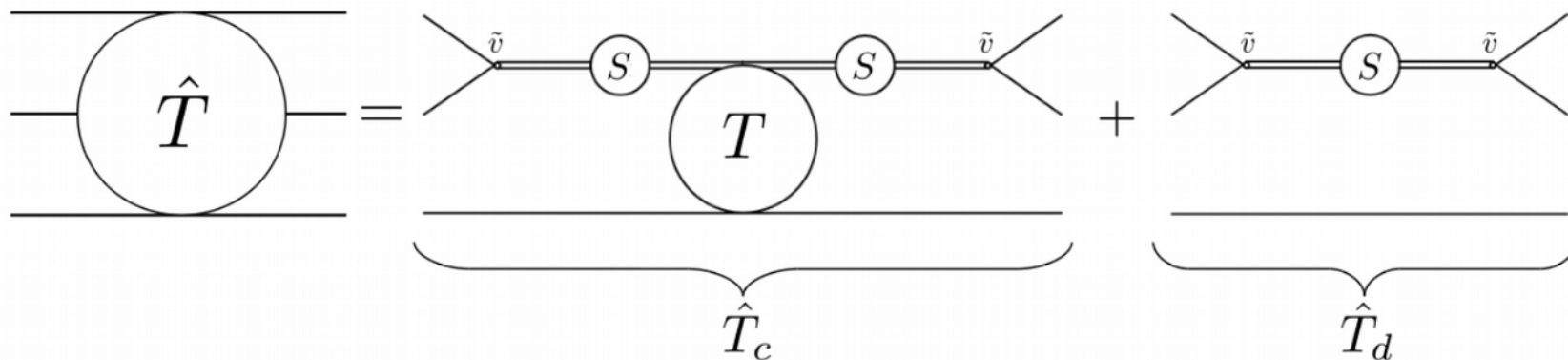
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Unitarity & Matching

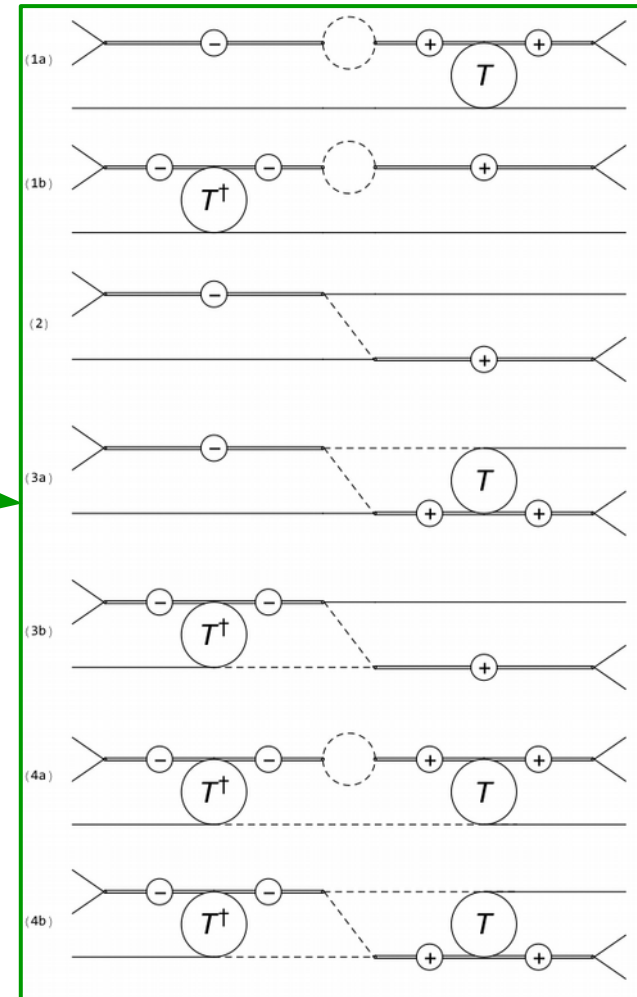
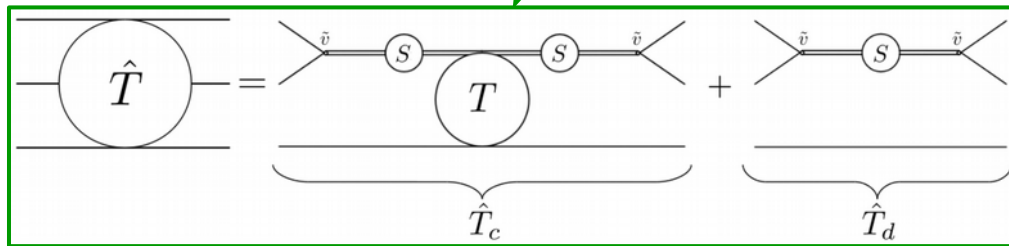
3-body Unitarity (normalization condition \leftrightarrow phase space integral)

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$

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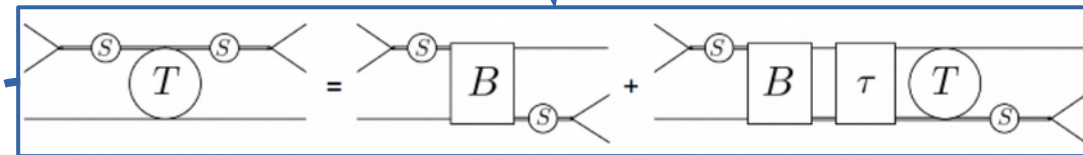
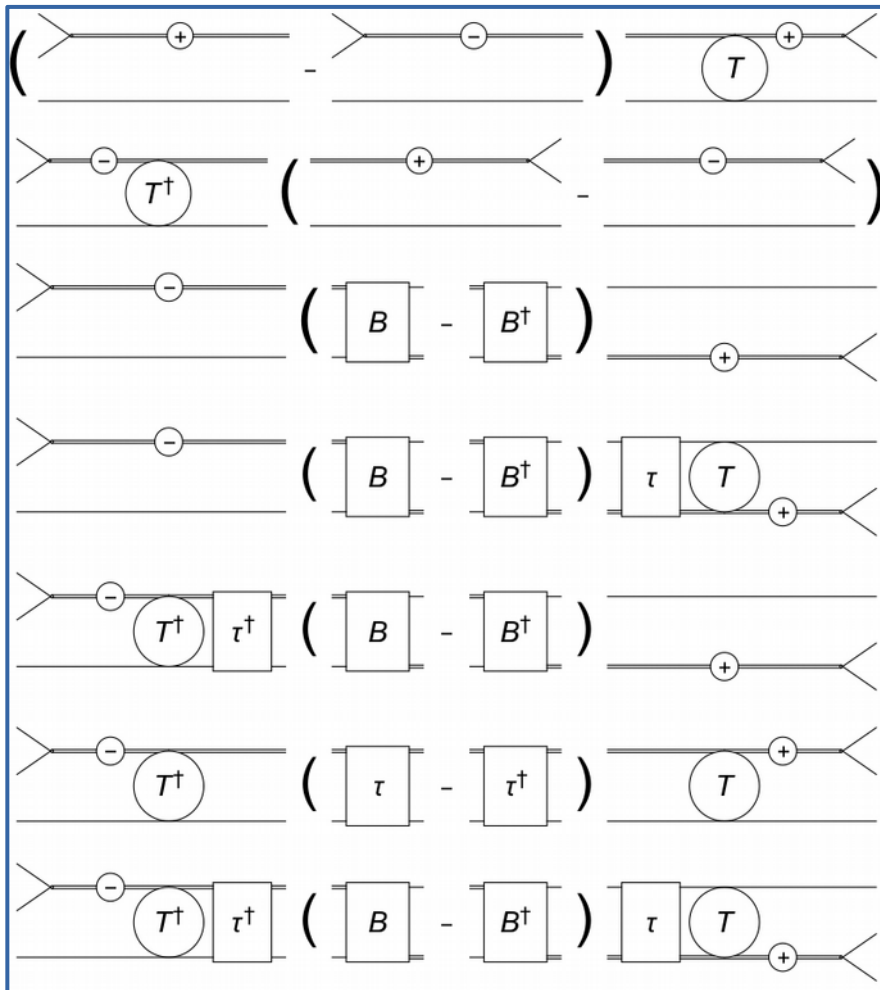
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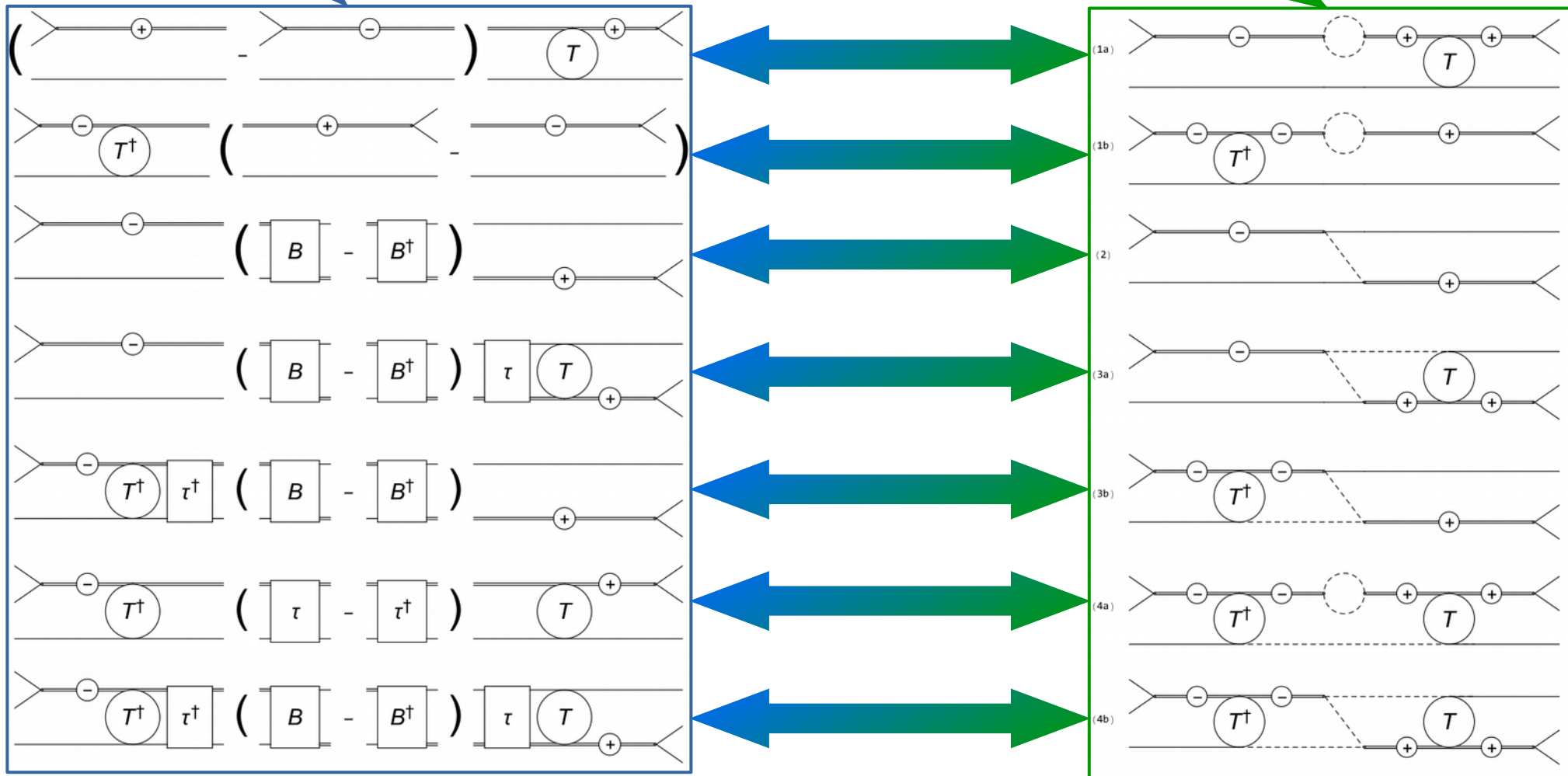


**General ansatz for the Isobar-spectator interaction
 $\rightarrow B$ & τ are unknown!!!**

Unitarity & Matching

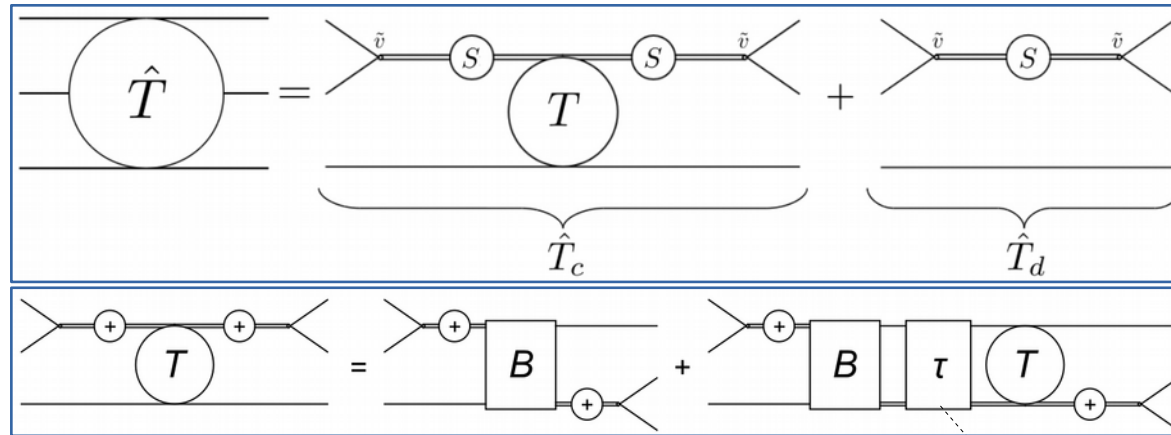
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SCATTERING AMPLITUDE

3 → 3 scattering amplitude is a 3-dimensional integral equation

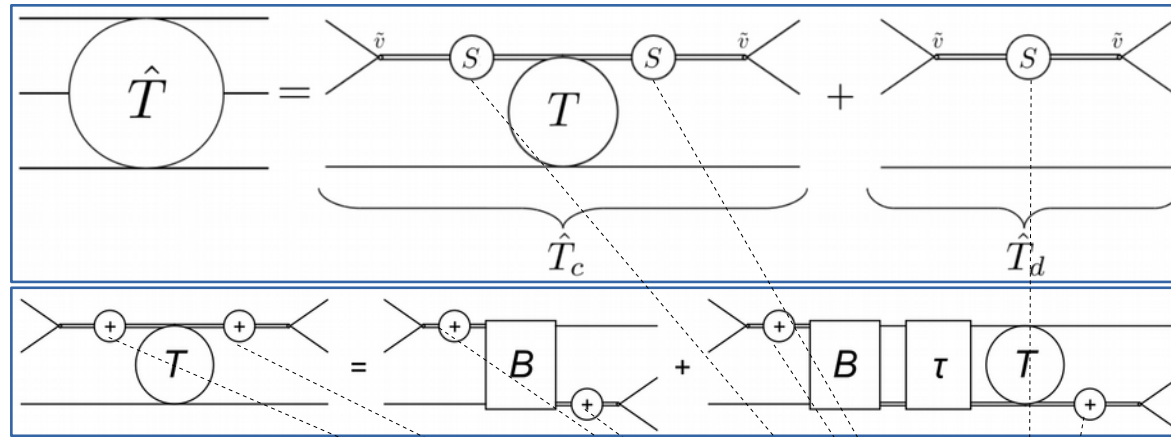


- Imaginary parts of B , S are fixed by **unitarity/matching**
- For simplicity $v=\lambda$ (full relations available)

$$\tau(\sigma(k)) = (2\pi)\delta^+(k^2 - m^2)S(\sigma(k))$$

SCATTERING AMPLITUDE

3 → 3 scattering amplitude is a 3-dimensional integral equation



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$$\text{Disc } \frac{1}{S} = -\frac{i}{8\pi} \frac{K_{\text{cm}}}{\sqrt{\sigma(k)}} \lambda^2$$

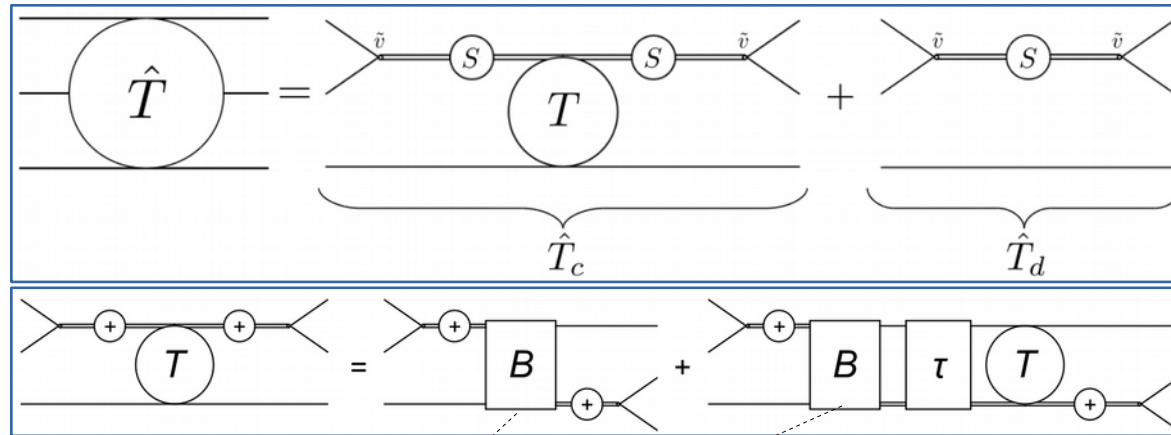
- twice subtracted dispersion relation in invariant mass - $\sigma(k)$

$$-\frac{1}{S} = \sigma(k) - M_0^2 - \frac{1}{(2\pi)^3} \int d^3\ell \frac{\lambda^2}{2E_\ell(\sigma(k) - 4E_\ell^2 + i\epsilon)}$$

- in the rest-frame of isobar (**Lorentz invariance!**)

SCATTERING AMPLITUDE

3 → 3 scattering amplitude is a 3-dimensional integral equation



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$$\text{Disc } B(u) = 2\pi i \lambda^2 \frac{\delta(E_Q - \sqrt{m^2 + Q^2})}{2\sqrt{m^2 + Q^2}}$$

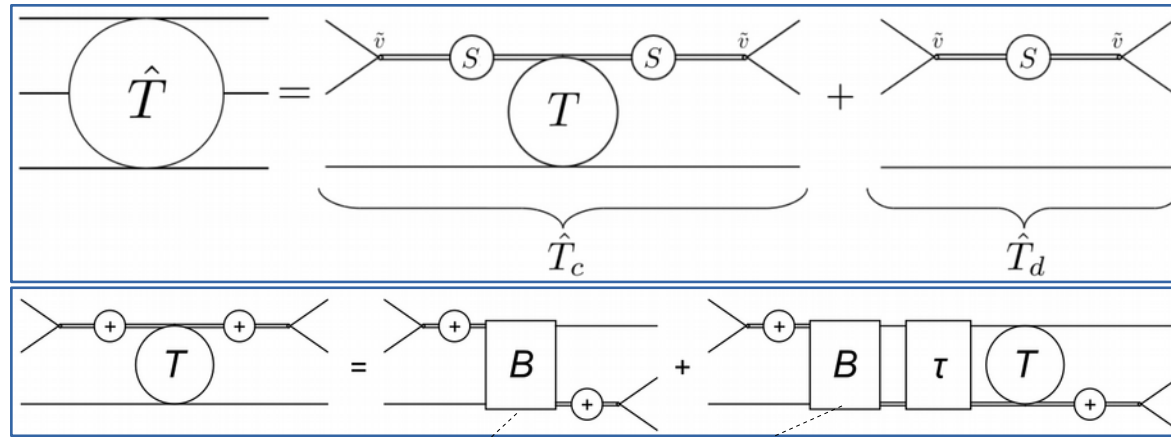
- un-subtracted dispersion relation

$$\langle q|B(s)|p\rangle = -\frac{\lambda^2}{2\sqrt{m^2 + Q^2} (E_Q - \sqrt{m^2 + Q^2} + i\epsilon)}$$

- one- π exchange in TOPT → **RESULT!**

SCATTERING AMPLITUDE

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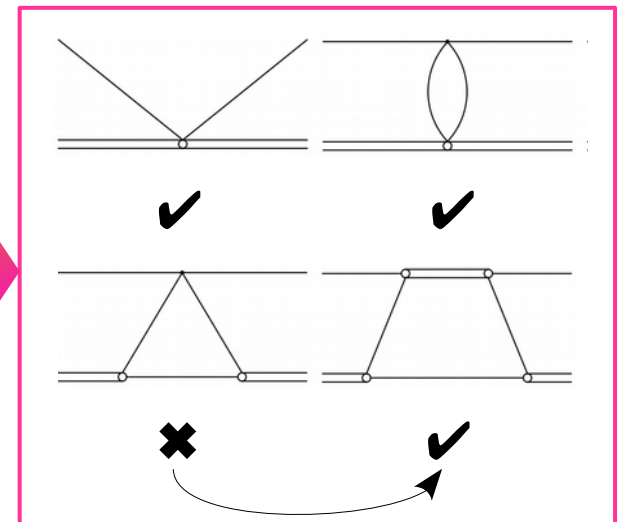
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THREE-BODY AMPLITUDE IN A BOX

MM, Döring

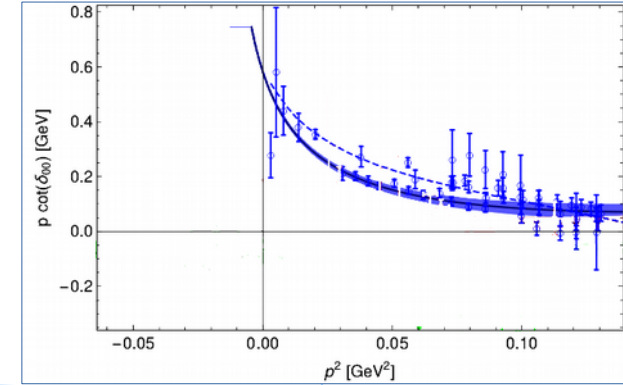
Arxiv: 1709.08222



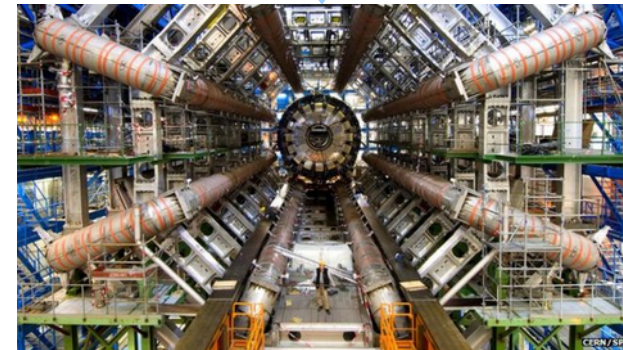
WHY LATTICE?

QCD

ab-initio



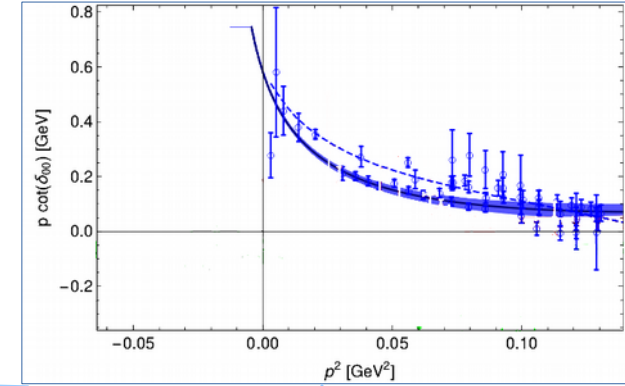
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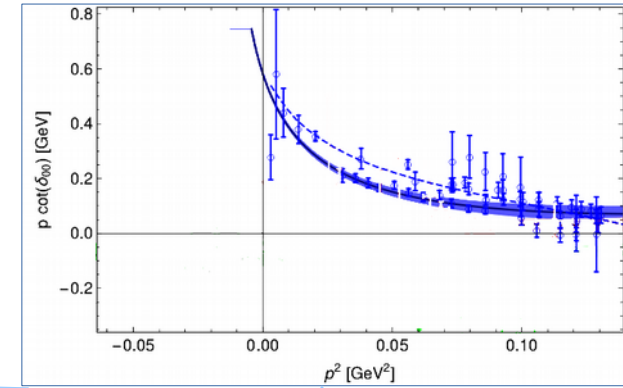
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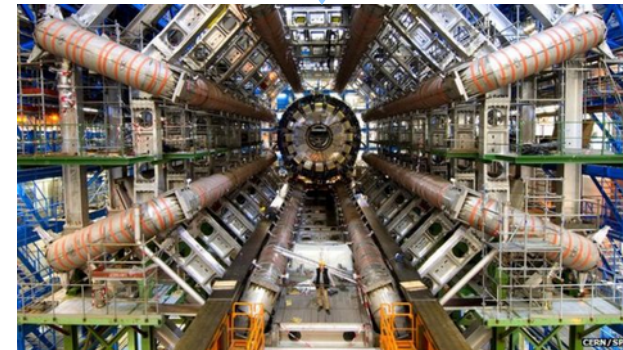
DATA



WHY LATTICE?



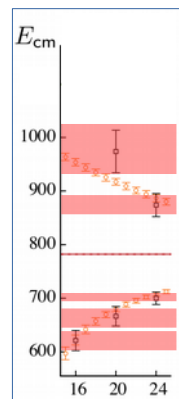
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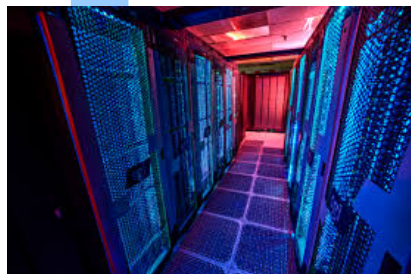
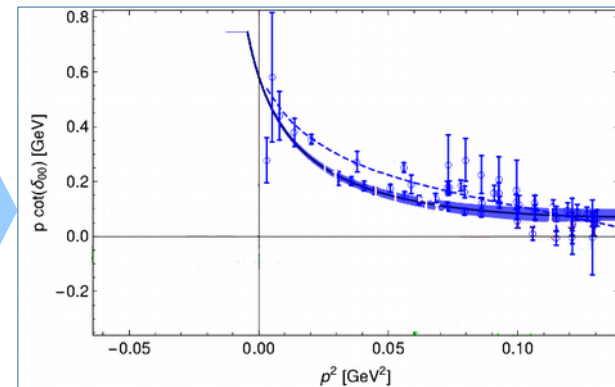
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Q C D

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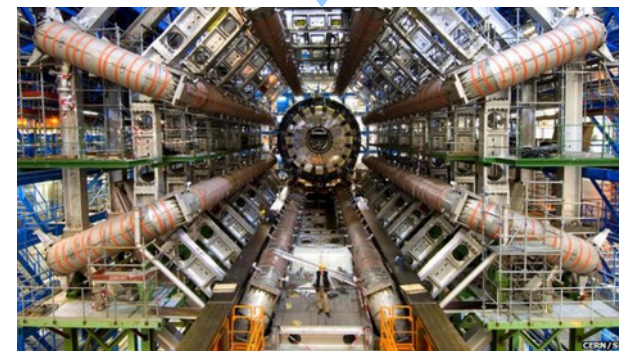


- 1) unphysical pion mass
- 2) (periodic) boundary conditions



~~ab-initio~~

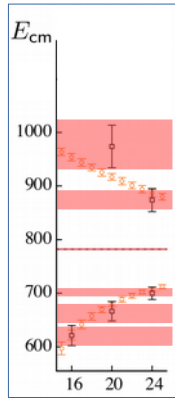
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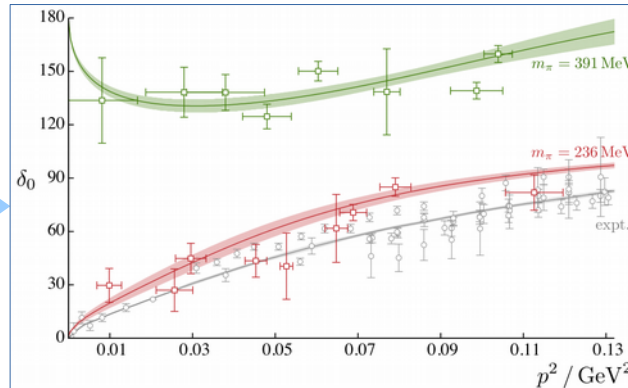
GOALS & CHALLENGES

Recipe for $2 \rightarrow 2$ scattering (e.g. $I=J=0$ $\pi\pi$ scattering)



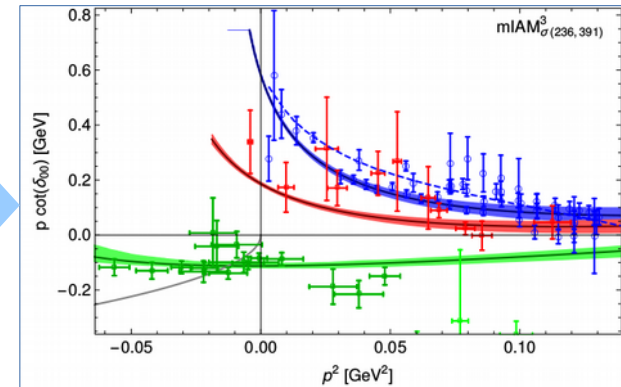
HSC(2016)

STEP 1



Briceño et al.(2016)

?



Doring, MM, Hu (2016)

LÜSCHER(1986)

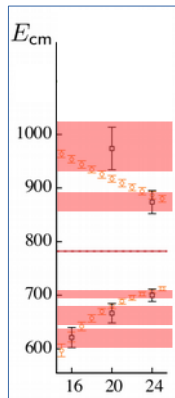
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He et al. (2005)

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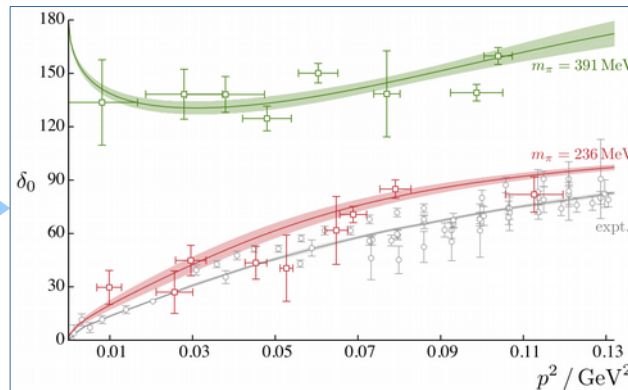
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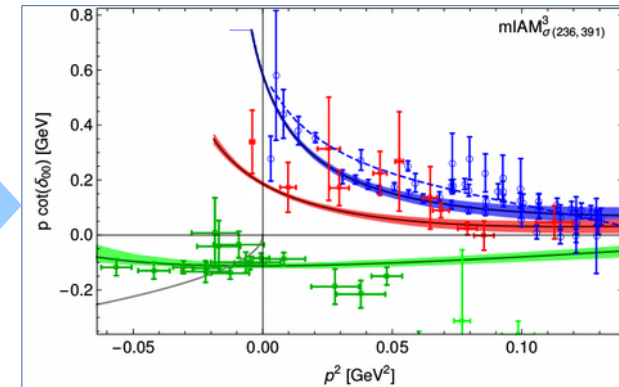
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CHIRAL EXTRAPOLATIONS

- M_π dependence from ChPT Gasser, Leutwyler(1981)
 - Extensions to resonances exist
- Hanhart et al. (2008)... Bruns, MM (2017)

GOALS & CHALLENGES

QCD calculations in finite volume

- 1) unphysical pion mass
- 2) (periodic) boundary conditions
→ **discrete momenta**

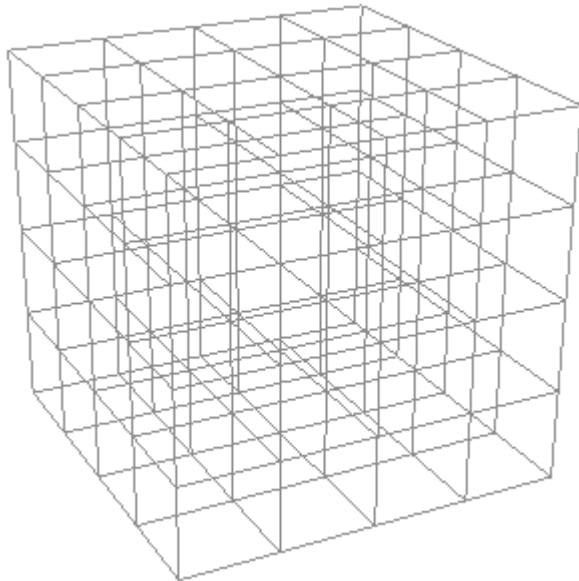
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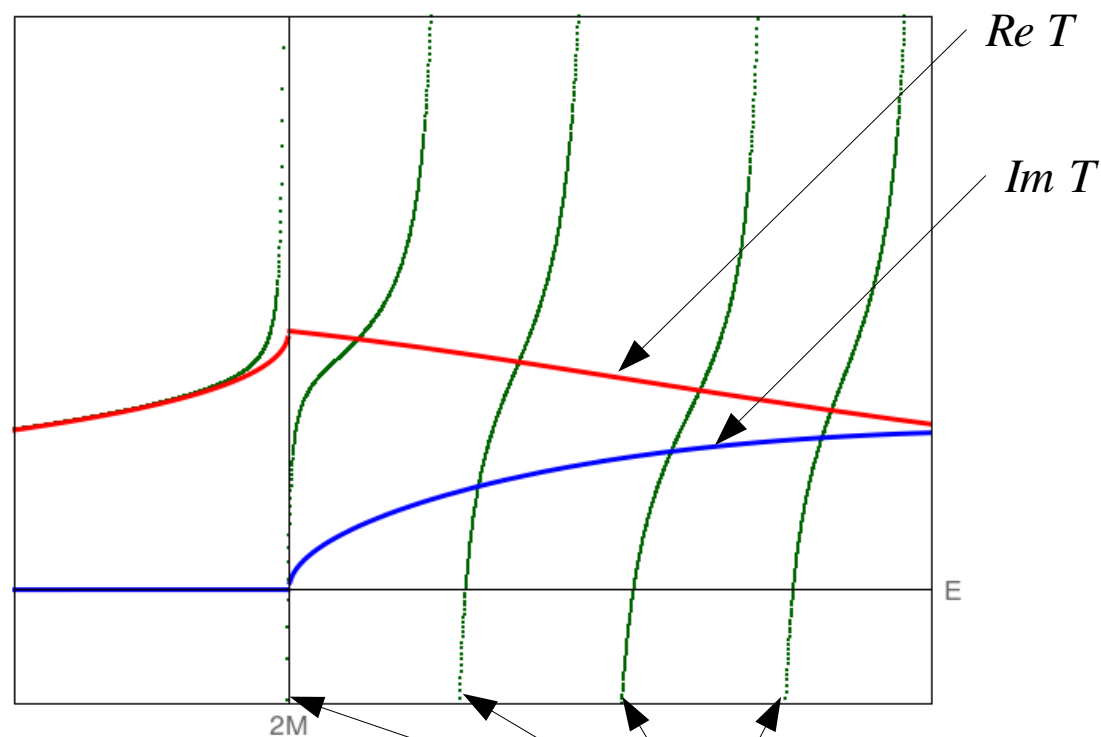
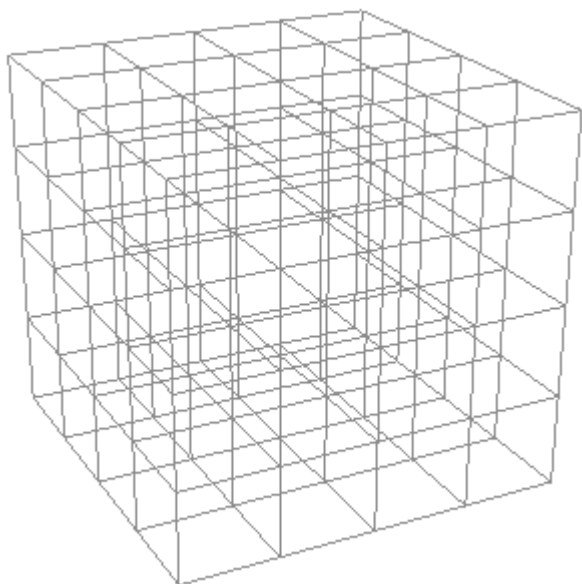
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Energy eigenvalues \leftrightarrow poles of T^{FV}

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Lüscher-like formalism in $3 \rightarrow 3$ case is under investigation

Polejaeva/Rusetsky (2012) Briceño/Hansen/Sharpe (2016)

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⇒ THIS WORK: discretize $3 \rightarrow 3$ scattering amplitude in isobar formulation

GOAL: quantization condition from 3-body unitarity!

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Partial Waves in infinite volume

- separation of angular momentum $\rightarrow Y_{lm}(\theta, \varphi)$
- reduces dimensionality of the problem

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& project to A_1^+ (basis vector: $Y_{00}(\theta, \varphi)$)*

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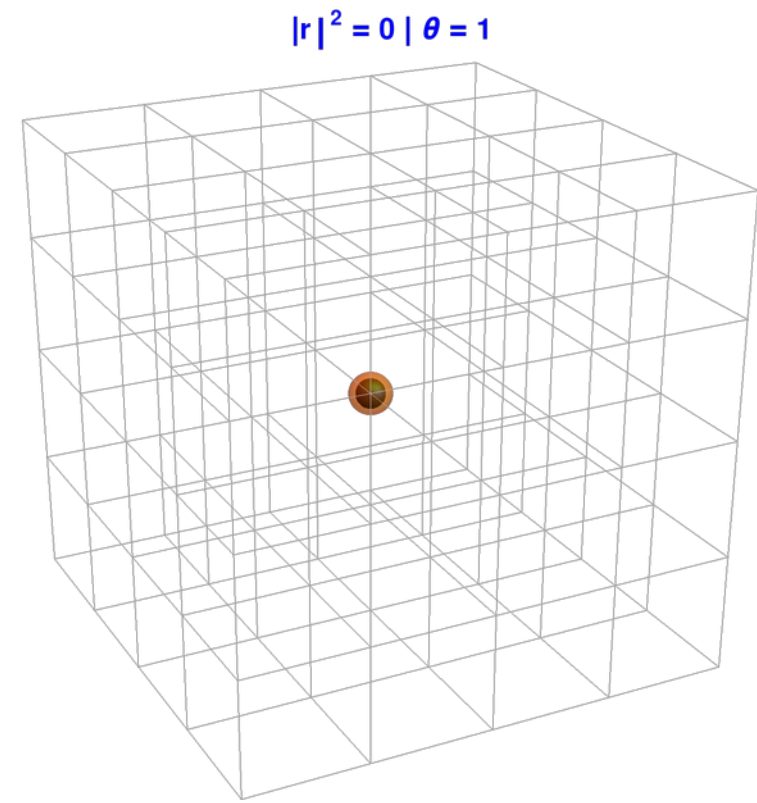


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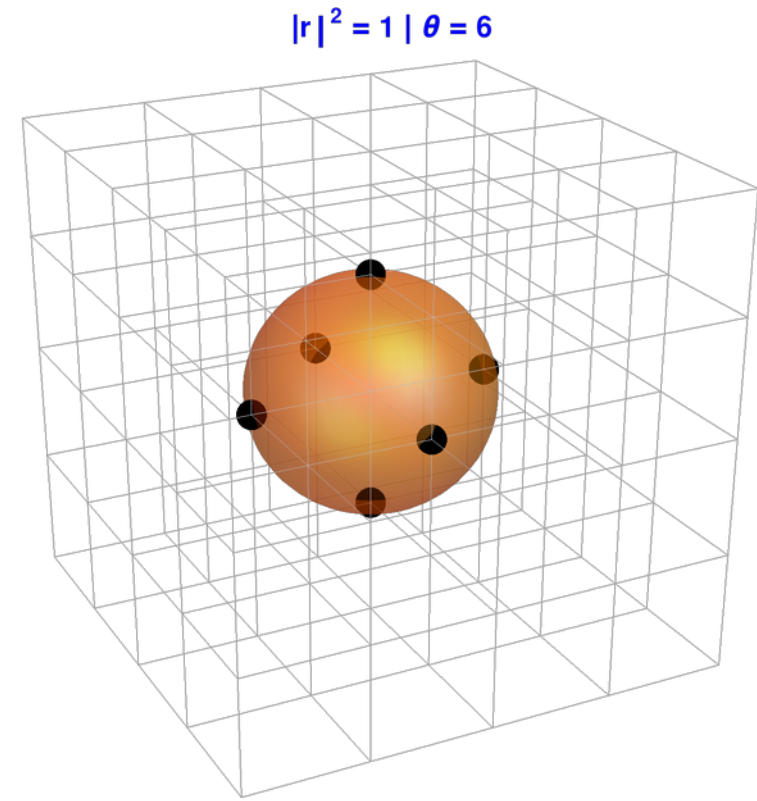


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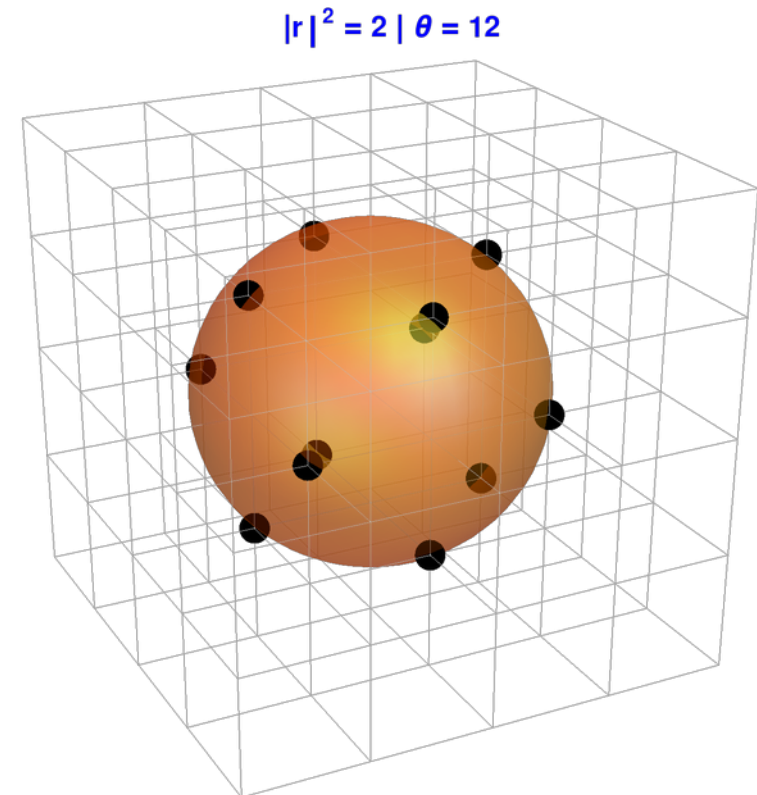


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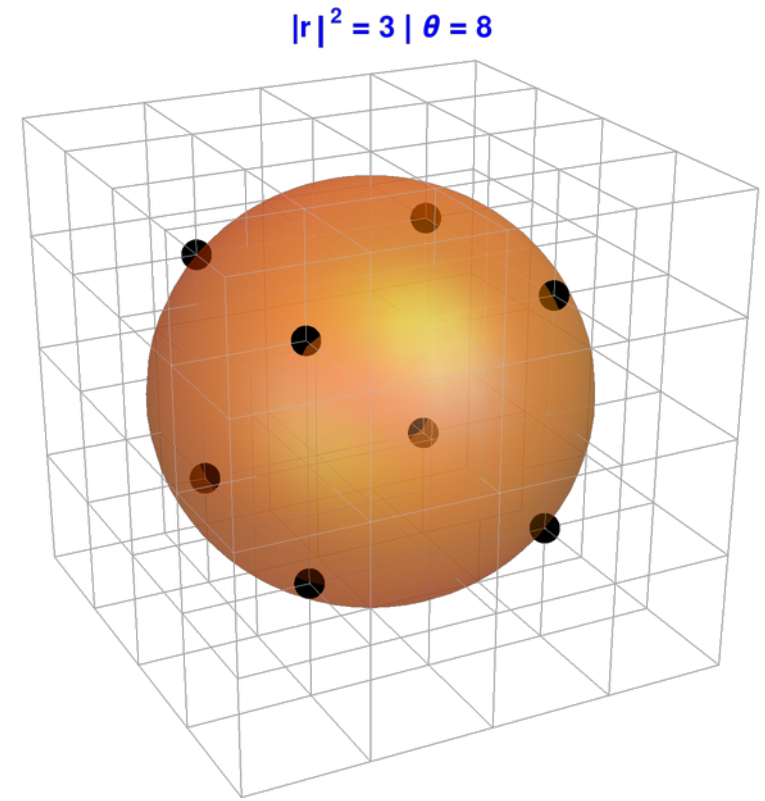


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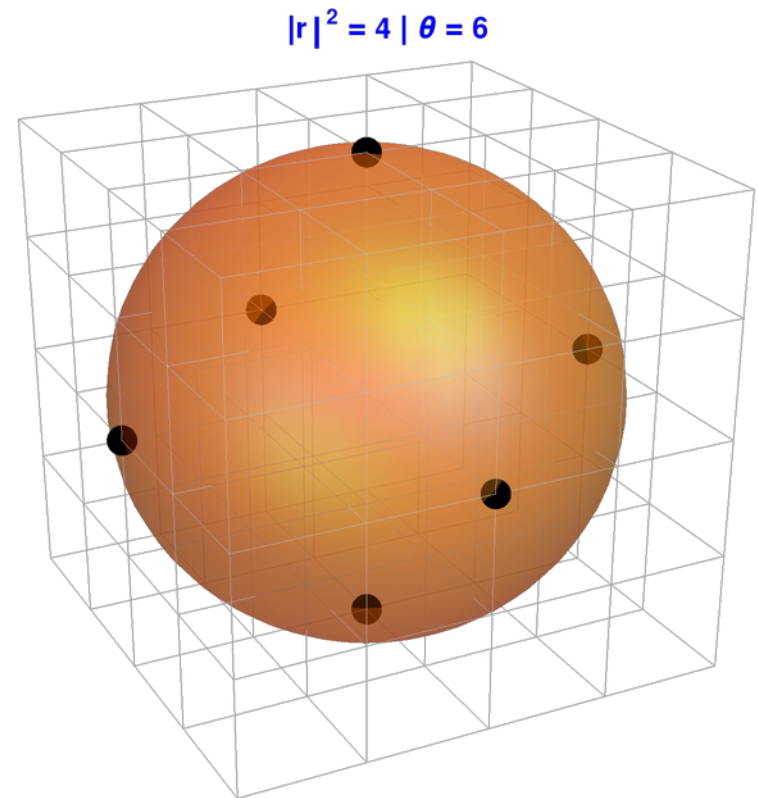


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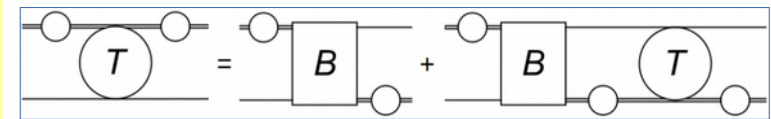
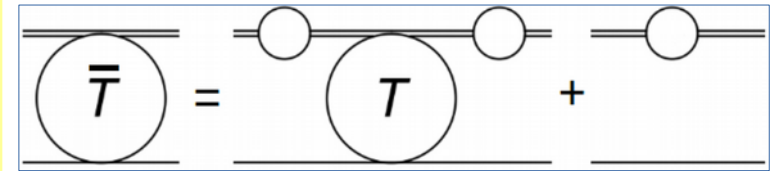
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$\bar{T}(W)$ is a matrix equation w.r.t $|q\rangle, |p\rangle=0,1,2,3,4,5,6,8$

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– when isobar-momenta are discretized in the 3-body cms momenta

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→ fin. vol. normalization of δ -distribution!

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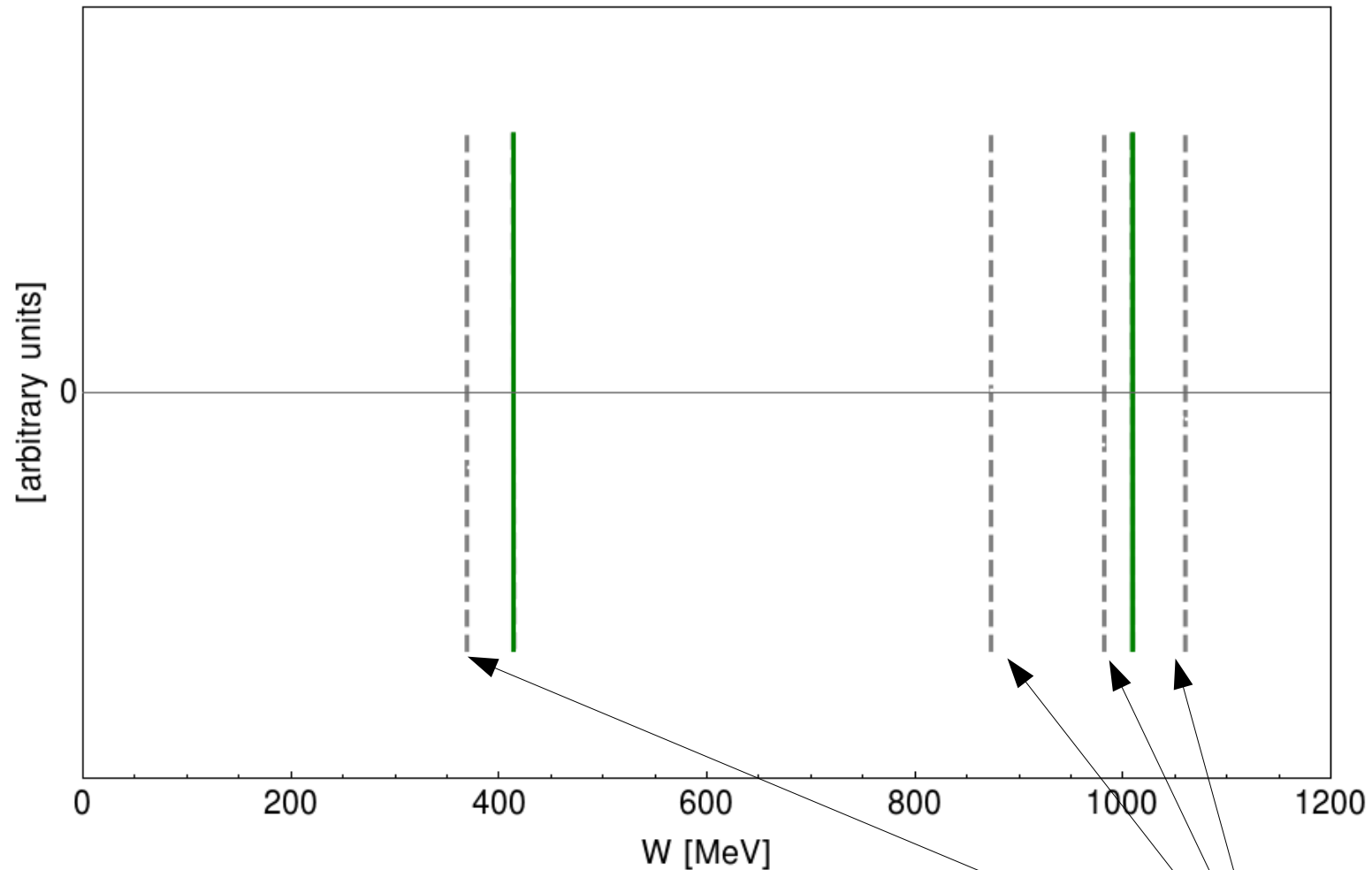
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Genuine 3-body eigenenergies = poles in s:

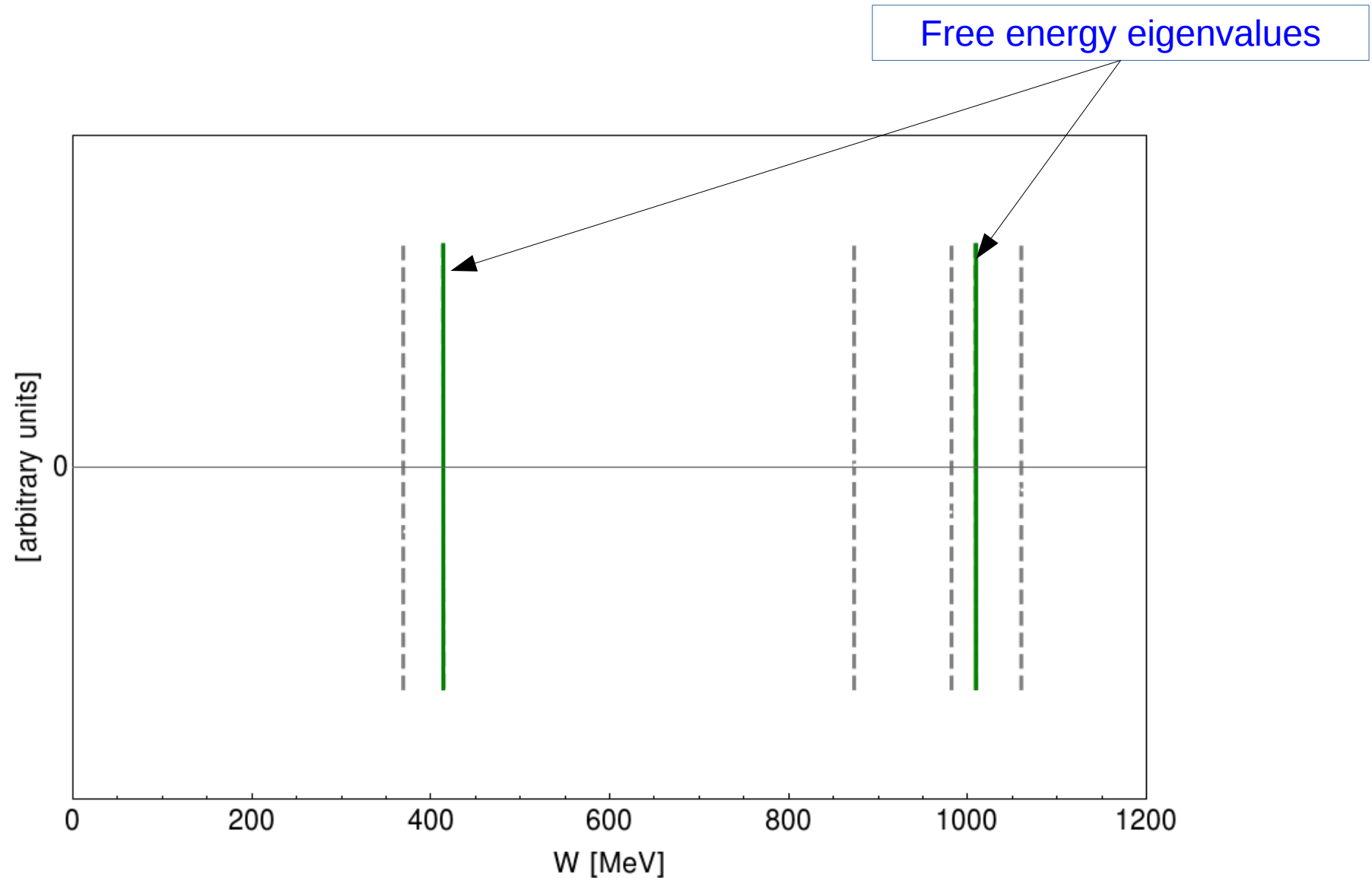
$$\text{Det} \left[B^{A_1^+}(s) \left[\frac{\vartheta(n)}{2E(s)L^3} \right] + \tau(s)^{-1} \right] = 0$$

RESULTS ($L=3$ fm, $M=138$ MeV)

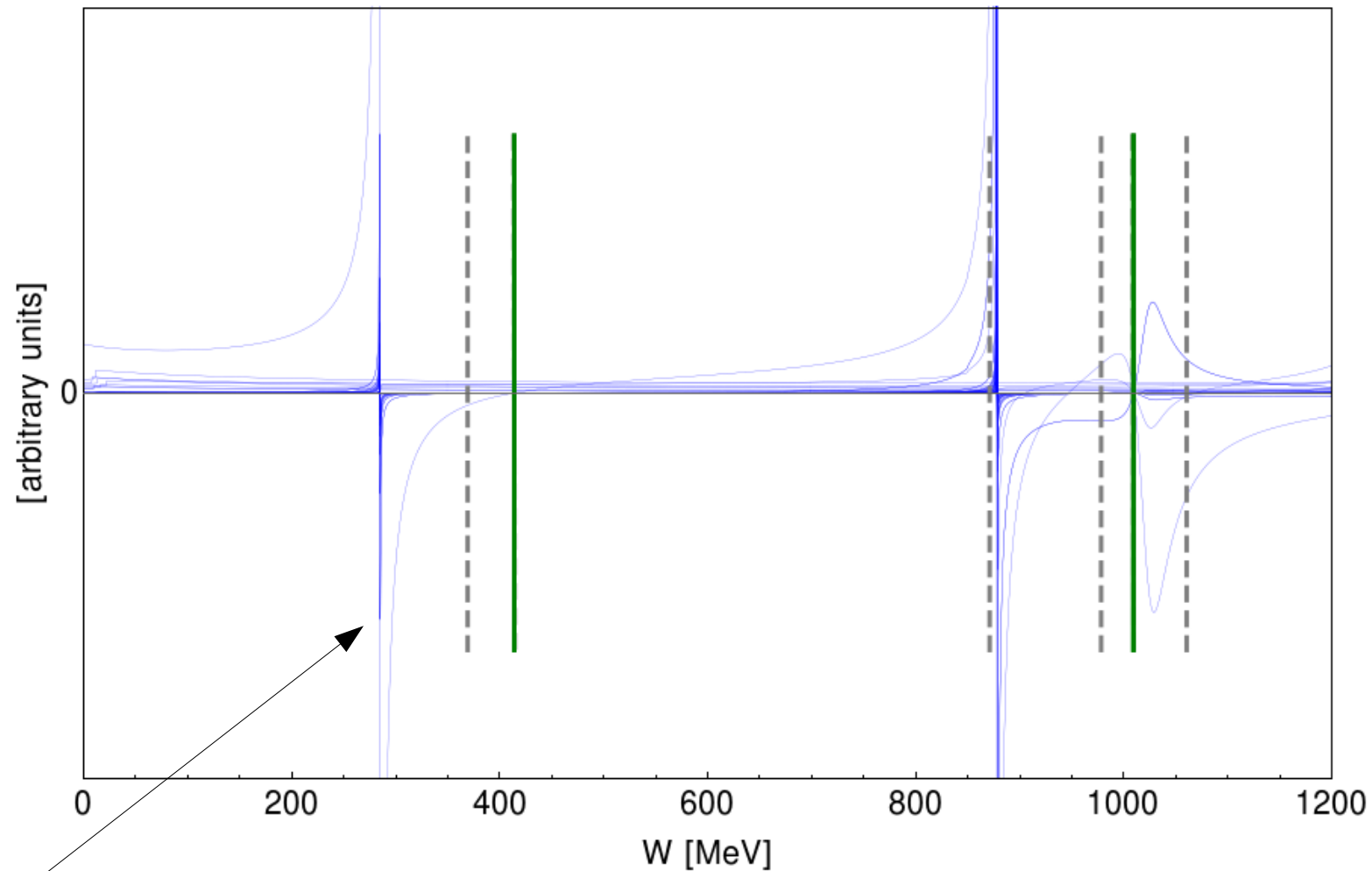


Isobar propagator poles

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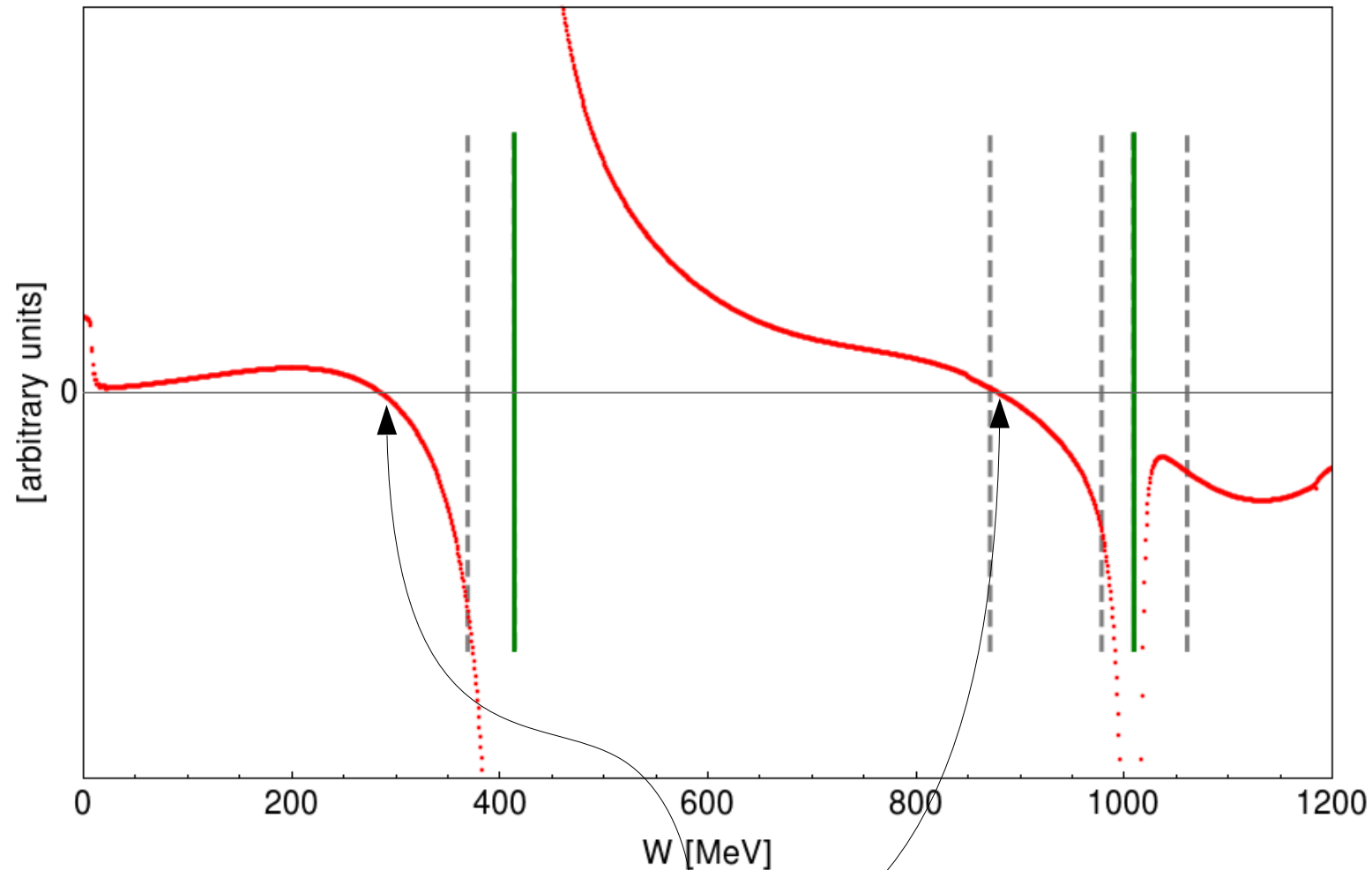


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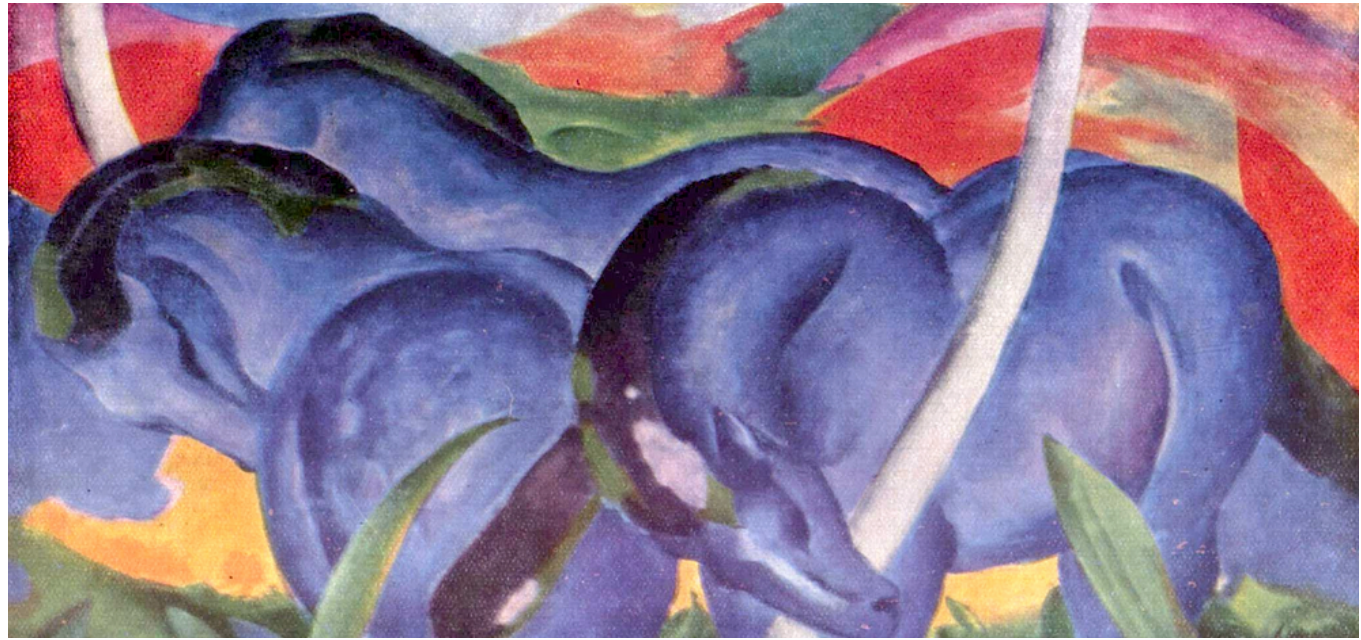
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SUMMARY



3-body amplitude in infinite volume

- 3-body Unitarity dictates imaginary parts of the driving term & isobar propagator
- Result: 3-dim. relativistic integral equations

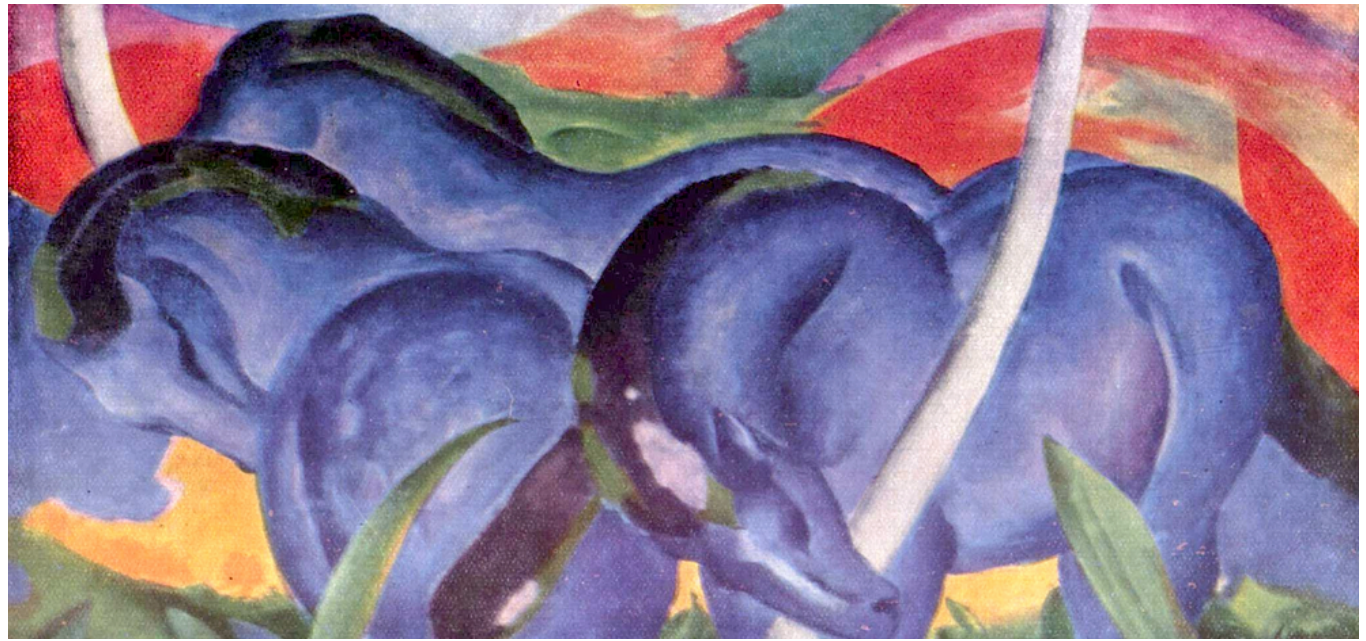
Finite volume investigation:

- Discretization techniques
- Quantization condition
- Case study → practicability!

OUTLOOK

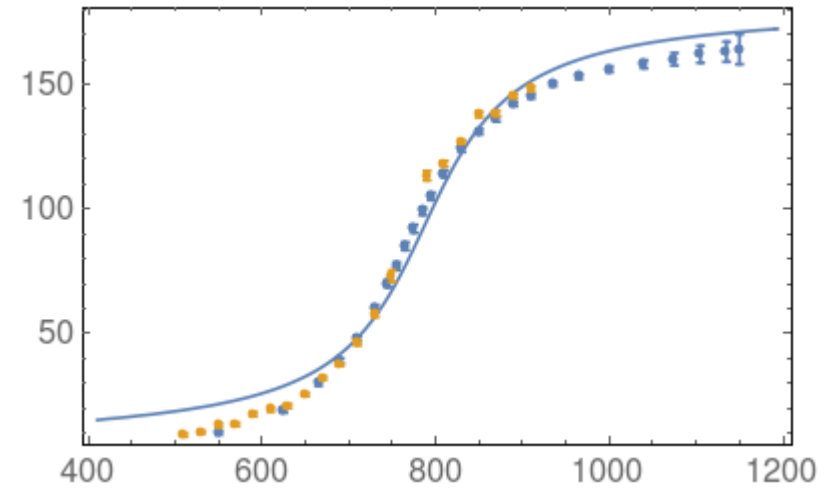
- include angular momentum / isospin / multiple isobars
- practical studies: $a_1(1260)$, ...

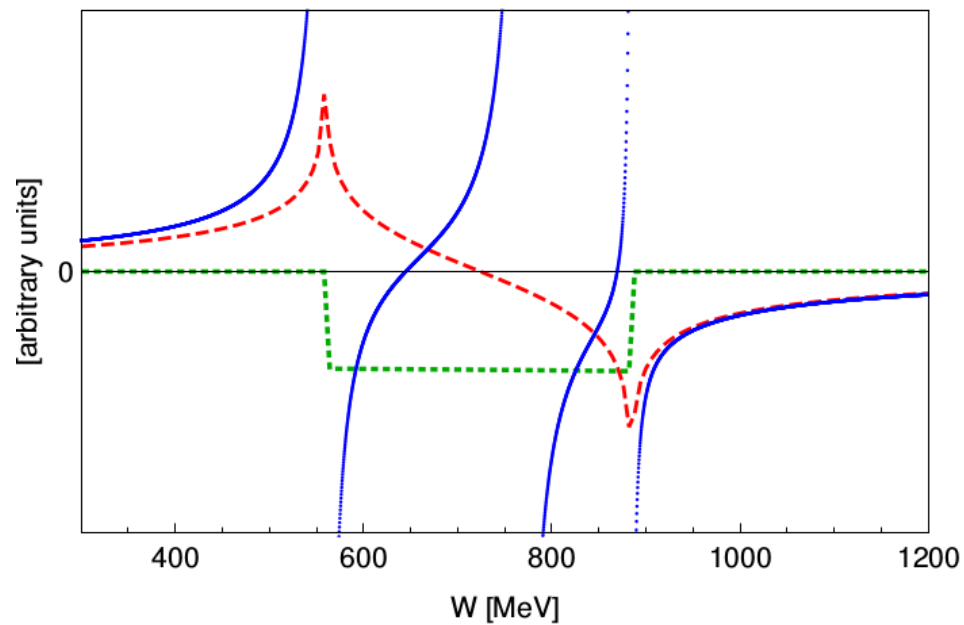
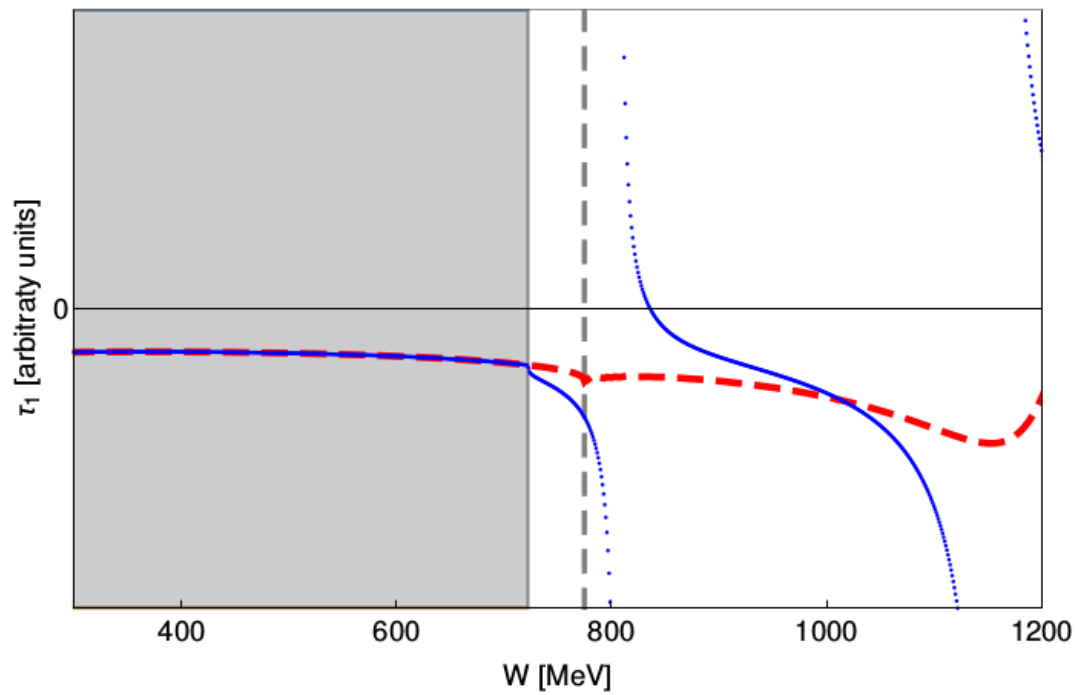
**THANK
YOU!**



SPARES

- $T_{22}(W) \approx v 1/D v$
- 3 free parameter: β (form factor), λ (strength of coupling), $M0$ (“bare mass of isobar”)
- Fixed to reproduce typical phase-shifts
 → just to get into the same ballpark





Unitarity & Matching

- 3-body Unitarity (normalization condition \leftrightarrow phase space integral)

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$

